

PROGRESS IN COMPUTATIONAL RHEOLOGY WITH THE FINITE VOLUME METHOD

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OUTLINE:

- Distinctive aspects of **Computational Rheology**
- Our activity in the field (1996-2005)
- **Constitutive equations** (differential)
- **FVM** (discretization)
 - Approximation of **convective fluxes**
 - **Stress/velocity** coupling
- Examples of results:
 - **Contraction** (*PTT* and *Old-B*) – *steady*
 - **Cylinder** (*Old-B* & *FENE*) – *steady/unsteady instability*
 - **Driven cavity** (*recoil: FENE*) – *unsteady*
 - **Expansion** (*FENE-CR*) – *steady*
 - **Vortex shedding** (*FENE*) – *unsteady*

Many others have made progresses in Computational Rheology with FVM:

- Tanner, Phan-Thien, Xue, and co-workers
- Webster and co-workers (mixed FEM and FVM)
- Owens, Sahin (FVM pressure-free)

But the talk concentrates on review of our work...

RELEVANCE OF COMPUTATIONAL RHEOLOGY:

PRACTICAL APPLICATIONS

- Industry of thermoplastics (PVC, etc)
- Extrusion and Molding processes
- Paints (non-dripping)
- Glues (super-glues...)
- Food stuffs (yogurths; spaghetti; ...)
- Cosmetics
- Two-phase flows
- Bioengineering: blood flow (hemodynamics)

DISTINGUISHING FEATURES OF COMPUTATIONAL RHEOLOGY :

- Elasticity and memory introduces a fluid time scale
- Level of elasticity measured by a Deborah number: $De = \lambda U / L$
- Additional equations for stresses (differential or integral)
- Stress equations are hyperbolic (full set, mixed type)
- Coupling between flow and constitutive eqs.
- Strength of coupling increases with “elasticity”: $\beta = \eta_s / \eta_0$
- Stiff coupling in time dependent flows (2 time scales)
- Real fluids require complicate constitutive eqs.
- Multimode (6 eqs x N modes, giving 30 eqs. for 5 modes)

SCOPE OF OUR WORK:

- Extend tools of classical CFD to Computational Rheology
- Finite Volume Method
- Original code developed at Imperial College (*Oliveira, 1992*)
(3D; non-orthogonal; non-staggered)
- Application to fundamental flows

ACHIEVEMENTS:

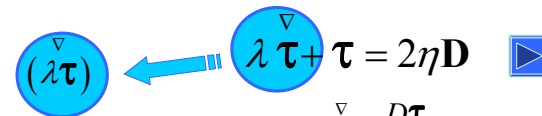
- General stress/strain solution code (many differential constitutive models)
- Method for stress/velocity coupling (stress at CV face)
- High Resolution Scheme (HRS) for stress equations
- Accurate results for some benchmark problems (contractions; expansions; cylinder; lid-driven cavity; etc)
- Ideas to improve the solution procedure (flow/stress linkage)

EQUATIONS:

Mass:
$$\frac{\partial \rho}{\partial t} + \nabla \cdot \rho \mathbf{u} = 0$$

Motion:
$$\rho \frac{D\mathbf{u}}{Dt} = -\nabla p + \nabla \cdot \boldsymbol{\tau} + \nabla \cdot [\eta_s (\nabla \mathbf{u} + \nabla \mathbf{u}^T)]$$

Constitutive (for the stress tensor $\boldsymbol{\tau}$):



$$\lambda \overset{\nabla}{\boldsymbol{\tau}} + \boldsymbol{\tau} = 2\eta \mathbf{D}$$

Oldroyd derivative:
$$\overset{\nabla}{\boldsymbol{\tau}} = \frac{D\boldsymbol{\tau}}{Dt} - (\boldsymbol{\tau} \cdot \nabla \mathbf{u} + \nabla \mathbf{u}^T \cdot \boldsymbol{\tau})$$

EXAMPLES OF CONSTITUTIVE MODELS (1):

Newtonian: $\boldsymbol{\tau} = 2\mu\mathbf{D} = \mu(\nabla\mathbf{u} + \nabla\mathbf{u}^T)$

Generalized Newtonian Fluid (GNF): $\boldsymbol{\tau} = \eta(\dot{\gamma})(\nabla\mathbf{u} + \nabla\mathbf{u}^T)$

viscosity model: $\eta(\dot{\gamma}) = K \dot{\gamma}^{n-1}$

Differential Viscoelastic:

$$\boldsymbol{\tau} + \lambda_{eq} \left(\frac{\partial \boldsymbol{\tau}}{\partial t} + \nabla \cdot (\mathbf{u}\boldsymbol{\tau}) \right) = \eta_{p,eq} (\nabla\mathbf{u} + \nabla\mathbf{u}^T) + \lambda_{eq} (\boldsymbol{\tau} \cdot \nabla\mathbf{u} + \nabla\mathbf{u}^T \cdot \boldsymbol{\tau})$$

$$\lambda_{eq} \equiv \lambda / f(\tau) \quad \eta_{p,eq} \equiv \eta_p / f(\tau)$$

EXAMPLES OF CONSTITUTIVE MODELS (2):

Oldroyd-B: $f(\tau) = 1 \Leftrightarrow \begin{cases} \lambda_{eq} \equiv \lambda \\ \eta_{p,eq} \equiv \eta_p \end{cases}$

UCM: Upper Convected Maxwell $\eta_s = 0$

FENE-CR: $\begin{cases} \lambda_{eq} \equiv \lambda / f(\tau) \\ \eta_{p,eq} \equiv \eta_p \end{cases}$

$$f(\tau) = \frac{L^2 + \frac{\lambda}{\eta_p} tr(\boldsymbol{\tau})}{L^2 - 3}$$

PTT Linear: $f(\tau) = 1 + \varepsilon \frac{\lambda}{\eta_p} tr(\boldsymbol{\tau})$

shear-thinning

PTT Exponential: $f(\tau) = \exp \left[\varepsilon \frac{\lambda}{\eta_p} tr(\boldsymbol{\tau}) \right]$

EXAMPLES OF CONSTITUTIVE MODELS (3):

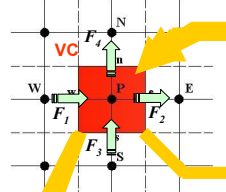
FENE-P:
$$f(\tau) = 1 + \frac{3}{b+2} \left[1 + \frac{\lambda}{3a\eta_p} \text{tr}(\tau) \right]$$

$$\begin{cases} \lambda_{eq} \equiv \lambda / f(\tau) \\ \eta_{p,eq} \equiv \eta_p / (f / a); a \equiv (b+5)/(b+2) \end{cases}$$

additional term:
$$\dots = \dots - \frac{D(1/f)}{Dt} (\lambda \tau + a \eta_p \mathbf{I})$$

Giesekus:
$$\begin{cases} \tau + \lambda \left(\tau + \frac{\alpha}{\eta_p} \tau \cdot \tau \right) = 2\eta_p \mathbf{D} \\ f(\tau) = 1 \end{cases} \quad \text{shear-thinning}$$

FINITE VOLUME METHOD (1):



$$\lambda \tau_{ij} + \left(\frac{\partial \tau_{ij}}{\partial t} + u_k \frac{\partial \tau_{ij}}{\partial x_k} \right) = \eta \left(\frac{\partial u_i}{\partial x_j} + \frac{\partial u_j}{\partial x_i} \right) + \left(\tau_{ik} \frac{\partial u_j}{\partial x_k} + \tau_{jk} \frac{\partial u_i}{\partial x_k} \right)$$

Mass conservation: (mass fluxes)

$$\sum_{f=1}^6 F_f = 0 \quad F_f = \sum_{j=1}^3 \rho B_{ff} \tilde{u}_j^{(n+1)}$$

Momentum conservation:

$$2^\circ \quad a_p u_p^{**} = \sum_{F=1}^6 a_F u_F^{**} + \left\{ -\nabla p^* + S_u [\nabla \cdot \tilde{\tau}^{(n+1)}] + S_u^{HOS} + \frac{\rho V}{\Delta t} (2.0 u_p^{(n)} - 0.5 u_p^{(n-1)}) \right\}$$

Stress evolution:

$$1^\circ \quad a_p \tau_p^{(n+1)} = \sum_{F=1}^6 a_F \tau_F^{(n+1)} + \left\{ S_\tau [\nabla u^*] + S_\tau^{HOS} + \frac{\lambda_{ef} V}{\Delta t} (2.0 \tau_p^{(n)} - 0.5 \tau_p^{(n-1)}) \right\}$$

Time advancement: 2nd order backward scheme

(Oliveira, Pinho, Pinto 1998; Oliveira 2001)

FINITE VOLUME METHOD (2):

Large sets of linearised equations $A.x=b$ solved with pre-conditioned conjugate gradient solvers.

Pressure-correction method (SIMPLEC).

$(p' = p^{**} - p^*)$ (akin to projection methods)

$$u_p^{(n+1)} = u_p^{**} - \frac{1}{a_p} \nabla p'$$
$$\nabla \cdot \left(\frac{1}{a_p} \nabla p' \right) - \{ \nabla \cdot u_p^{**} \} = 0$$

Non-linearities dealt with by iteration.
(inside a time step, for time dependent calculations)

FINITE VOLUME METHOD (3):

Numerical Issues:

A – Stress/velocity coupling.

B – Discretization of convective fluxes.

A - STRESS/VELOCITY COUPLING METHOD:

Term needed in u -eq.: $\nabla \cdot \tau \rightarrow \sum_{f=1}^6 \sum_{j=1}^3 B_{fj} (\tilde{\tau}_{ij})_f = \sum_{f=1}^6 T_{fi}$

Write stress eq.:as:

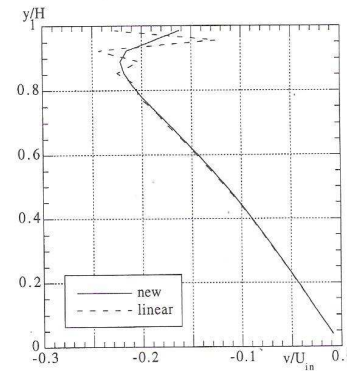
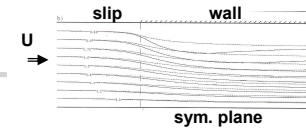
$$a_p^\tau \tilde{\tau}_{ij}^p = \sum_F a_F^\tau \tilde{\tau}_{ij}^{(n-1)} + \sum_{l=1}^3 \left(b_{li} [\Delta u_j]_l + b_{lj} [\Delta u_i]_l + \frac{b_{ll}}{\Delta t} [\Delta u_l]_l \right) + \frac{S_{ij}^\tau}{\Delta t} (2.0 \tilde{\tau}_{ij}^{(n)} - 0.5 \tilde{\tau}_{ij}^{(n-1)})$$

where: $b_{li} = \eta B_{li} + \lambda \sum_{k=1}^3 B_{lk} \tau_{ik}$ and B_{li} are areas

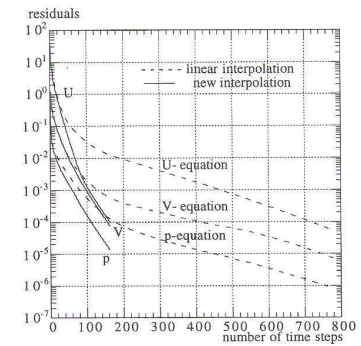
Divide by central coeff. a_p^τ and interpolate:

$$\tilde{\tau}_{ij} = \sum_{l \neq f} a'_F \tilde{\tau}_{ij}^F + \sum_{l=1}^3 \left(b'_{li} [\Delta u_j]_l + b'_{lj} [\Delta u_i]_l \right) + \tilde{b}'_{fi} [\Delta u_j]_f + \tilde{b}'_{fj} [\Delta u_i]_f + \tilde{S}_{ij}^\tau$$

EXAMPLE:



Vertical profile of v-velocity in entry flow



Residuals of equations

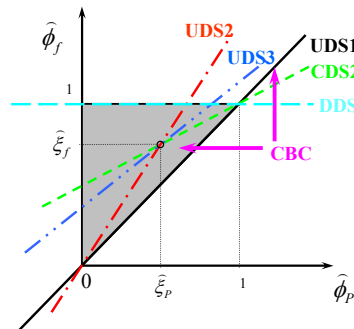
Slip/Stick Flow

B - CUBISTA: HIGH RESOLUTION SCHEME (1):

Calculation of cell-face flux:

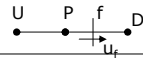
$$F_f \phi_f$$

- Important for accuracy
- Upwind (UDS1) too inaccurate
- High Order (HOS): oscillations
- Avoid oscillations: HRS
- Problem: iterative convergence



(B.P. Leonard)

HOS in Normalised Variable Diagram (NVD): $\hat{\phi} = \frac{\phi - \phi_U}{\phi_D - \phi_U}$



B - CUBISTA: HIGH RESOLUTION SCHEME (2):

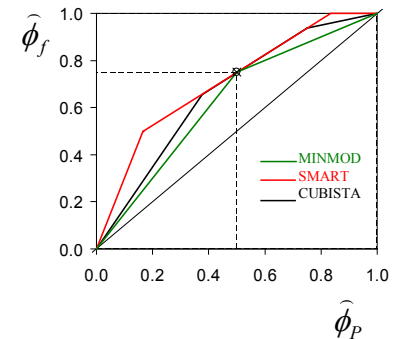
Definition of scheme:

$$\hat{\phi}_f = \begin{cases} \frac{7}{4} \hat{\phi}_p & 0 < \hat{\phi}_p < \frac{3}{8} \\ \frac{3}{4} \hat{\phi}_p + \frac{3}{8} & \frac{3}{8} \leq \hat{\phi}_p \leq \frac{3}{4} \\ \frac{1}{4} \hat{\phi}_p + \frac{3}{4} & \frac{3}{4} < \hat{\phi}_p < 1 \\ \hat{\phi}_p & \text{elsewhere} \end{cases}$$

(based on TVD restrictions)

$$\hat{\phi}_f \leq (2 - C) \hat{\phi}_p; C = 0$$

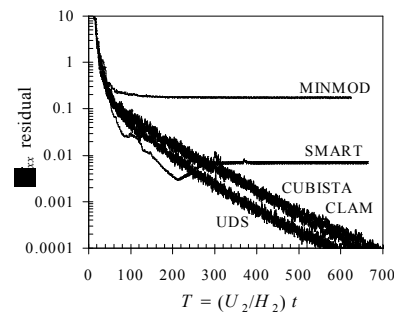
we chose: $C = 0.25$



(Alves, Oliveira, Pinho, IJNMF 2003)

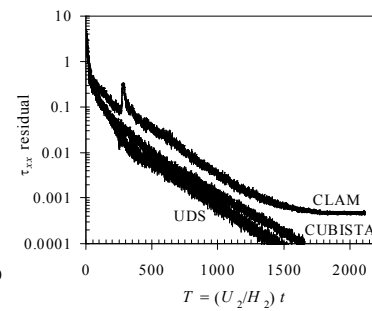
B - CUBISTA: HIGH RESOLUTION SCHEME (3):

Contraction flow, UCM, $De=3$



Mesh M2

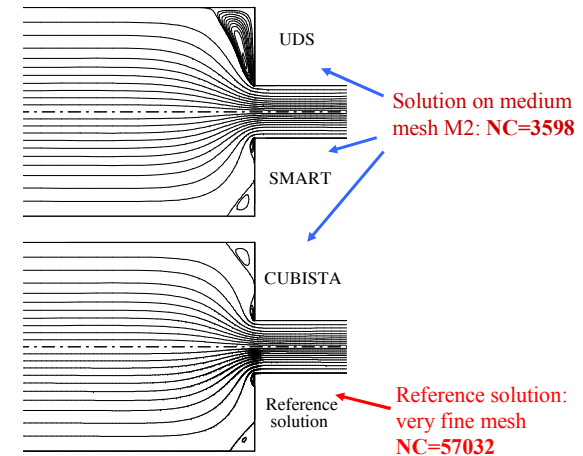
(NC=3598, $Dx=0.02$)



Mesh M3

(NC=14258, $Dx=0.01$)

B - CUBISTA: HIGH RESOLUTION SCHEME (4):

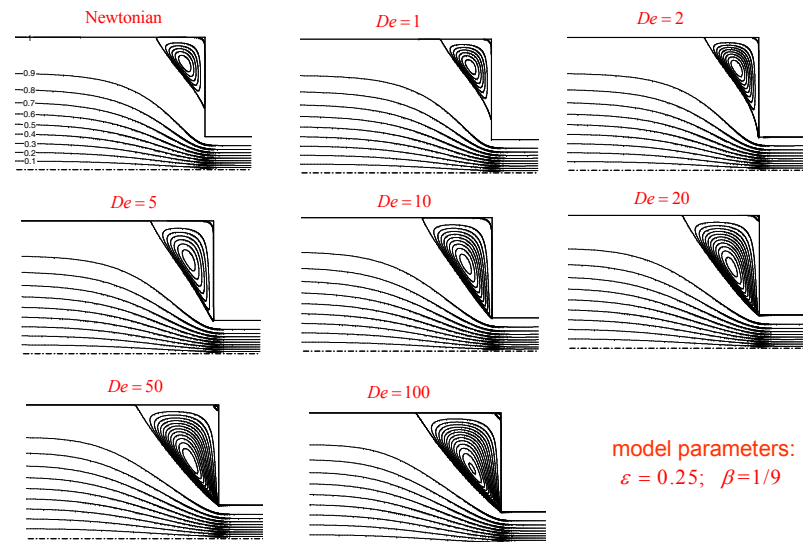


EXAMPLES OF CALCULATIONS:

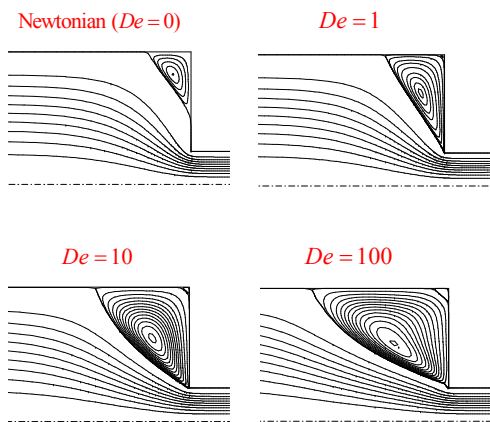
- Flow through contractions
- Flow around confined cylinder
- Lid-Driven cavity flows
- Flow through expansions
- Vortex shedding behind unbound cylinder
- Pulsating pressure gradient



4:1 PLANAR CONTRACTION FLOW, PTT LINEAR

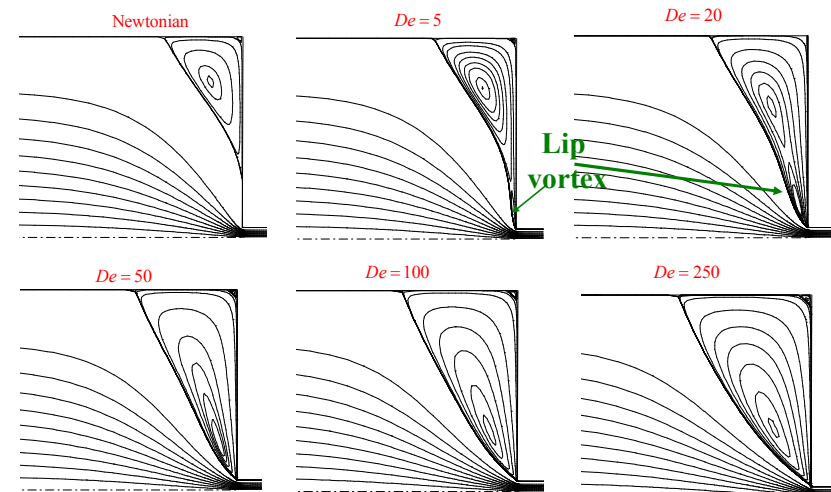


4:1 AXISYMMETRIC CONTRACTION FLOW, PTT LINEAR



(Alves, Pinho, Oliveira, ENCIT 2004)

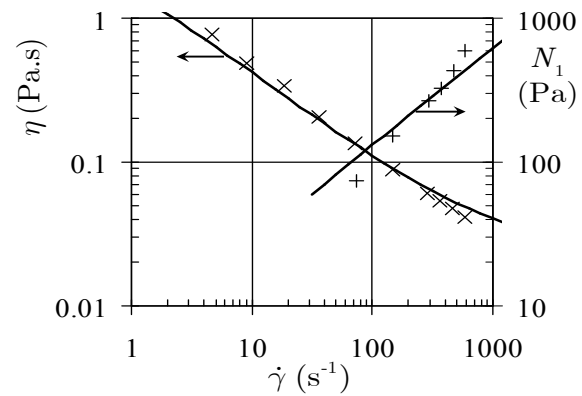
20:1 PLANAR CONTRACTION FLOW, PTT LINEAR



(Alves, Oliveira, Pinho, JNNFM 2004)

Comparison with experiments, 16:1 (Evans & Walters 1986)

1% polyacrylamide in water fitted with a PTT

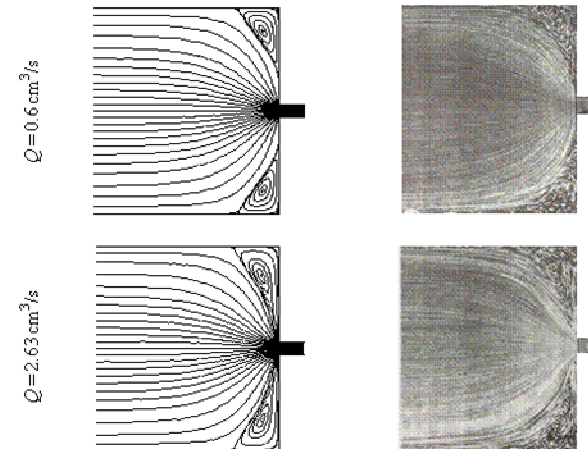


$\epsilon = 0.25$; $\beta = 1/120$
 $\lambda = 2.0$ s; $\eta_0 = 2.5$ Pa.s

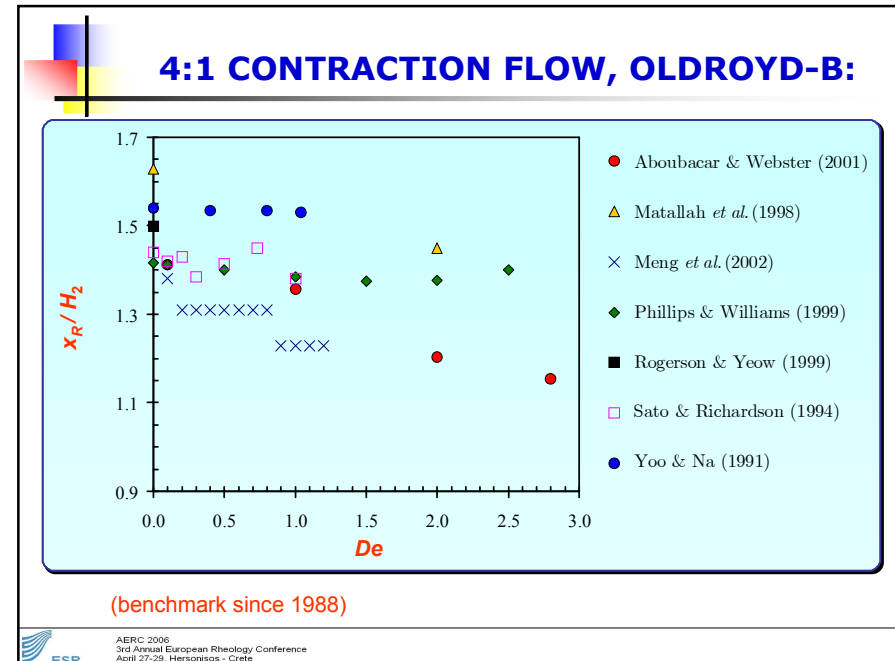
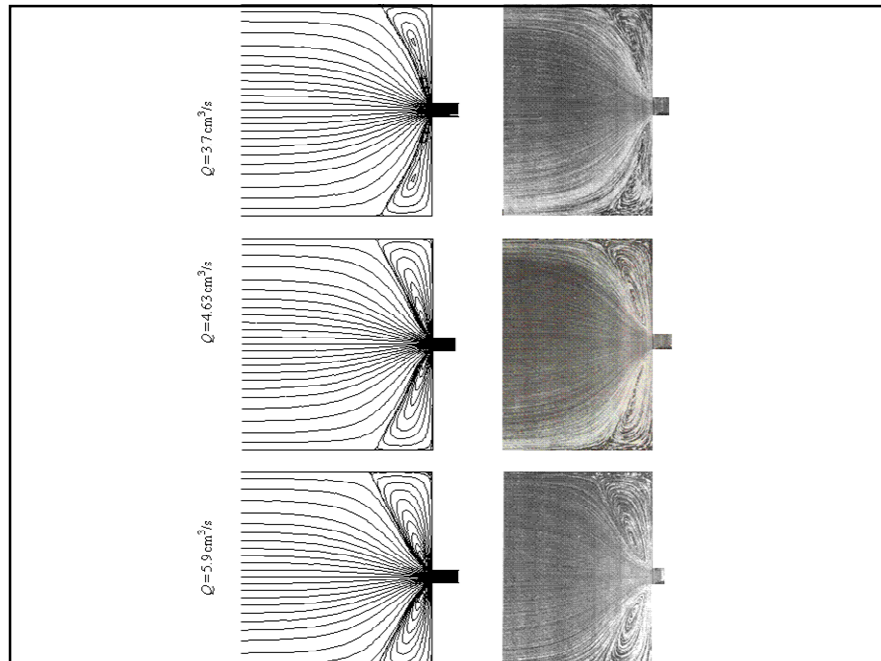


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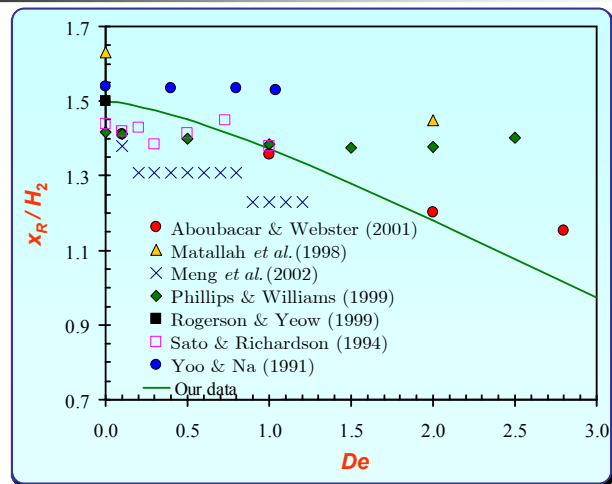
Comparison with experiments, 16:1 (Evans & Walters 1986)



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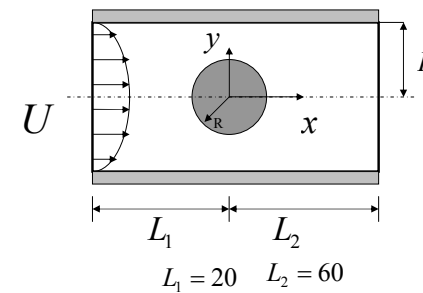


4:1 CONTRACTION FLOW, OLDROYD-B:



Alves, Oliveira, Pinho, JNNFM (2000; 2003)

THE CONFINED CYLINDER BENCHMARK:



“blockage”:

$$B = R / H = 0.5$$

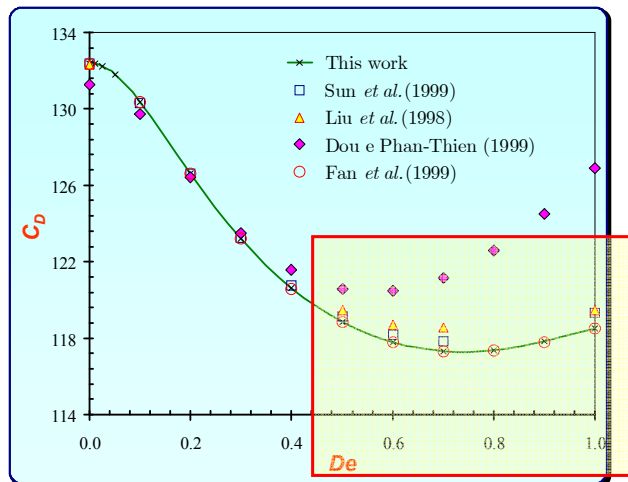
$$De = \lambda U / R$$

$$Re = 0$$

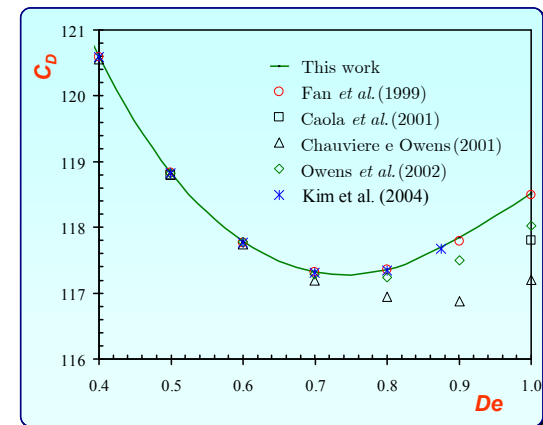
(benchmark since 1994)

CONFINED CYLINDER WITH OLDROYD-B:

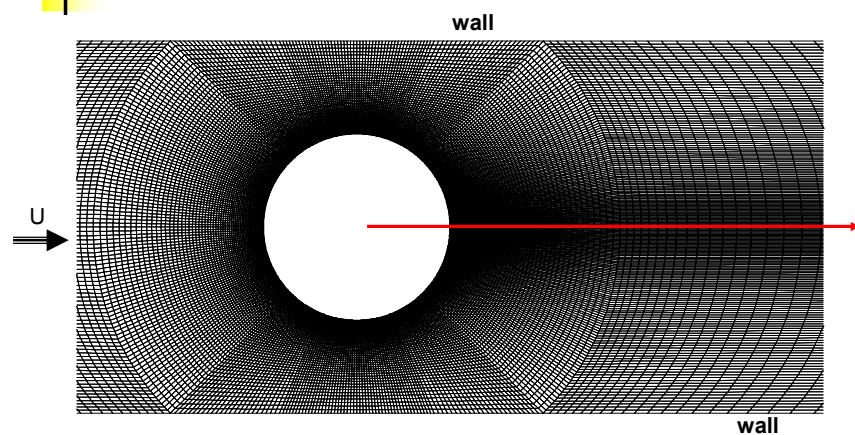
Initial calculations: *Alves, Pinho, Oliveira, JNNFM (2001)*



CONFINED CYLINDER WITH OLDROYD-B:



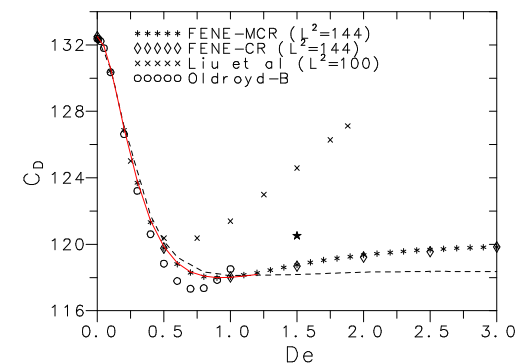
FENE-CR CALCULATIONS: example of mesh



Zoom in $-3 < x < 5$; M60(WR)-FO

- refined along wake
- full domain
- row of cells along wake

STEADY FLOW: DRAG COEFFICIENT (FENE-CR)



L^2
extensibility
parameter

Red: $L^2=144$
Dash: $L^2=100$
(finer mesh)

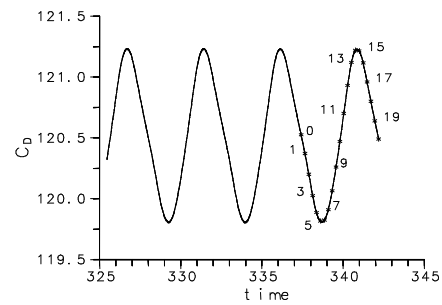
Sahin, Owens



(Oliveira & Miranda, JNNFM 2005)

UNSTEADY REGIME – DRAG vs. TIME:

Natural Instability (Hopf)



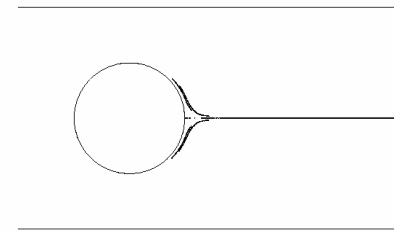
$$\Delta t = 2.4 \times 10^{-3}$$

$$T \approx 4.7$$

$$De=1.5; Re=0; L^2=144; \beta=0.59$$

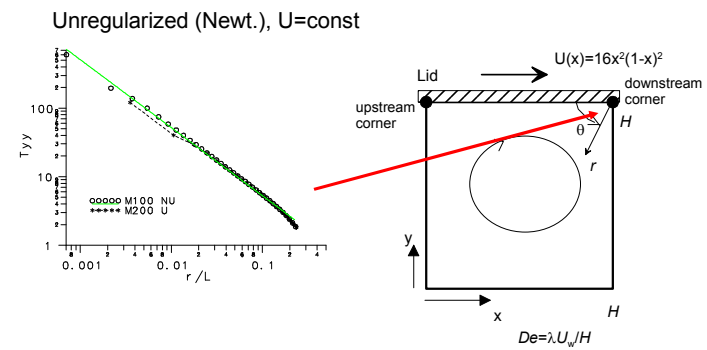


EXAMPLE OF STREAKLINES, DE=1.5:



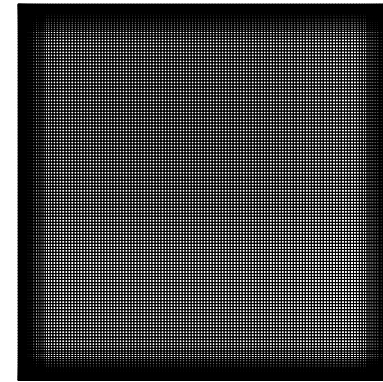
Note: no undulation

LID-DRIVEN CAVITY:



(Oliveira, IWNMNNFM 2005)

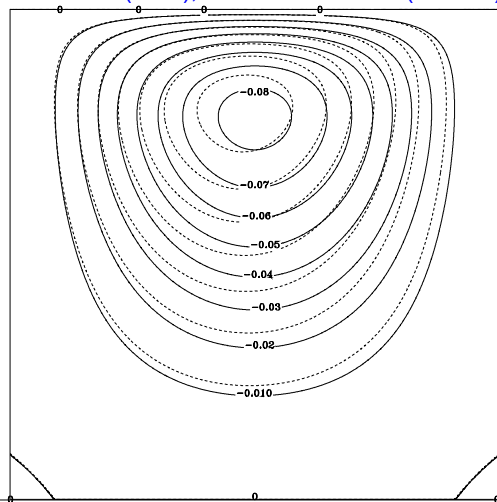
MESH-4: 160x160 NU



$$\Delta x_{\min} = 1.0 \times 10^{-3}$$

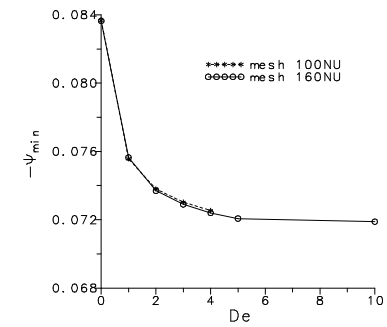
STREAMLINE COMPARISON:

Newtonian (solid), Viscoelastic $De=1$ (dashed)

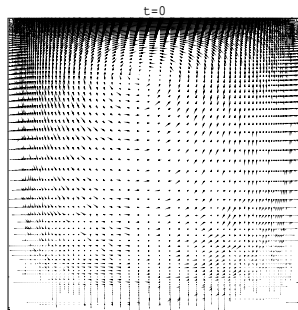


Pakdel et al.
PoF 1998

VARIATION OF MAIN EDDY RECIRCULATION WITH ELASTICITY:



TIME DEPENDENT FLOW: RECOIL FROM $De=2$



N1-steady



Velocity
Vectors



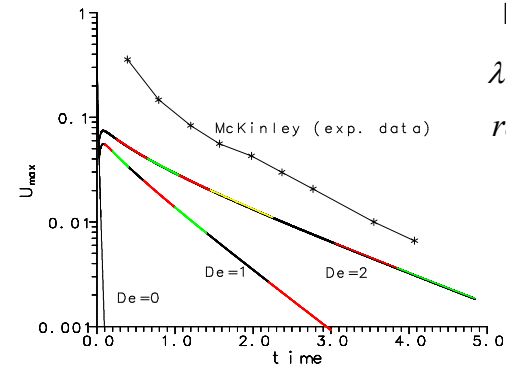
Streamlines



$De=2$; Mesh 4

$\Delta t = 0.5 \times 10^{-4}$
up to:
 $\Delta t = 5 \times 10^{-4}$

DECAY OF MAXIMUM VELOCITY IN CAVITY:



McKinley:

$\lambda = 1.6s$; $rate = 1.13s$

$rate / \lambda = 0.71$

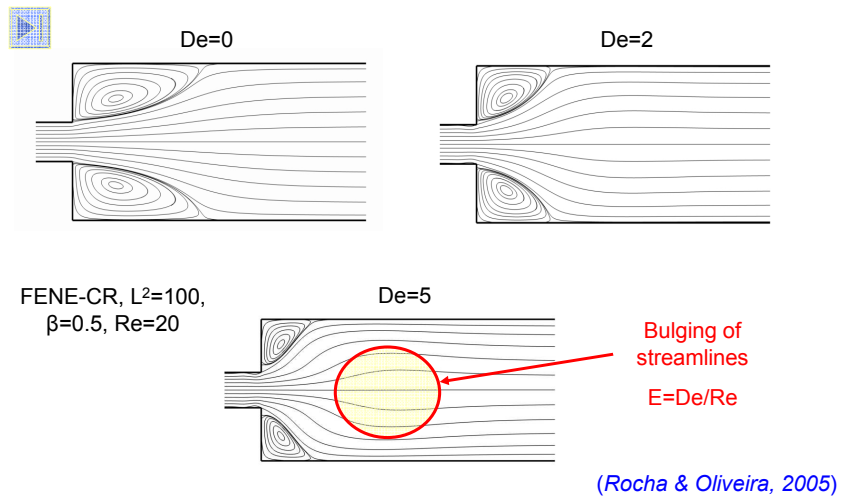
$De=0, T=0.019$

$De=1, T=0.75$

$De=2, T=1.48$

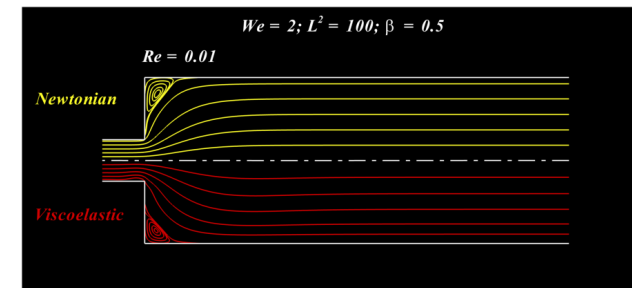
$T/De=0.75; 0.74$

FLOW THROUGH EXPANSIONS (1:4):



1:4 PLANAR EXPANSION:

$Re \leq 35$, no bifurcation



EFFECT OF POLYMER CONCENTRATION (β):

$$C = (1/\beta) - 1$$

Newtonian case:

$$Re = 20; We = 0; L^2 = 100$$

$$\beta = 1.0$$

Viscoelastic case 1:

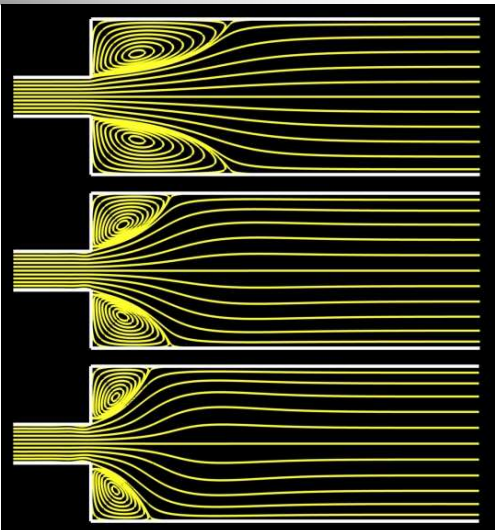
$$Re = 20; We = 2; L^2 = 100$$

$$\beta = 0.5$$

Viscoelastic case 2:

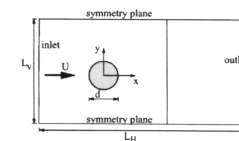
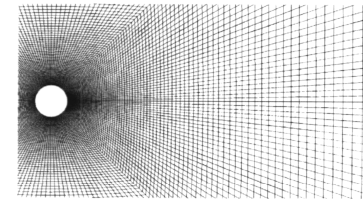
$$Re = 20; We = 2; L^2 = 100$$

$$\beta = 0.3$$



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VORTEX SHEDDING BEHIND UNCONFINED CYLINDER: mesh & flow domain



$$Re = \frac{\rho U d}{\eta_0} = 100$$

$$De = \frac{\lambda U}{d} = 0.80$$

$$L^2 = 100$$

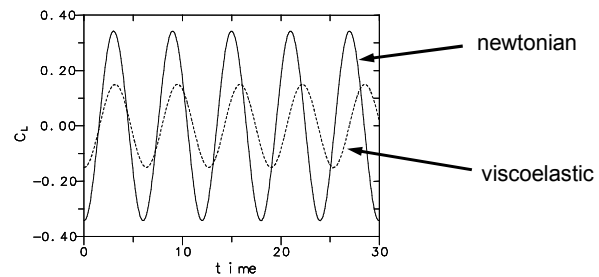
$$c = 0.1; \beta = 0.909$$

FENE-CR : constant-viscosity elastic liquid



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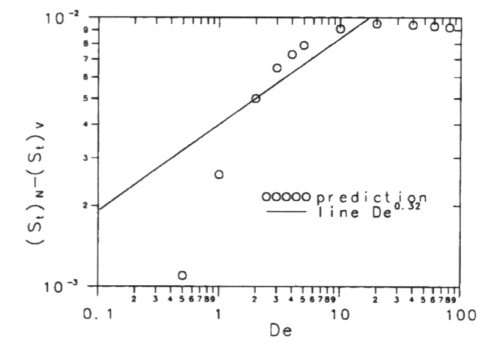
VORTEX SHEDDING BEHIND UNCONFINED CYLINDER: effect on lift



Lift coefficient versus time
 $De=80$, $L^2=100$, $Re=100$

- period: +5.8%
- magn. lift: -56 %
- magn. drag: -83%

VORTEX SHEDDING BEHIND UNCONFINED CYLINDER: effect on frequency

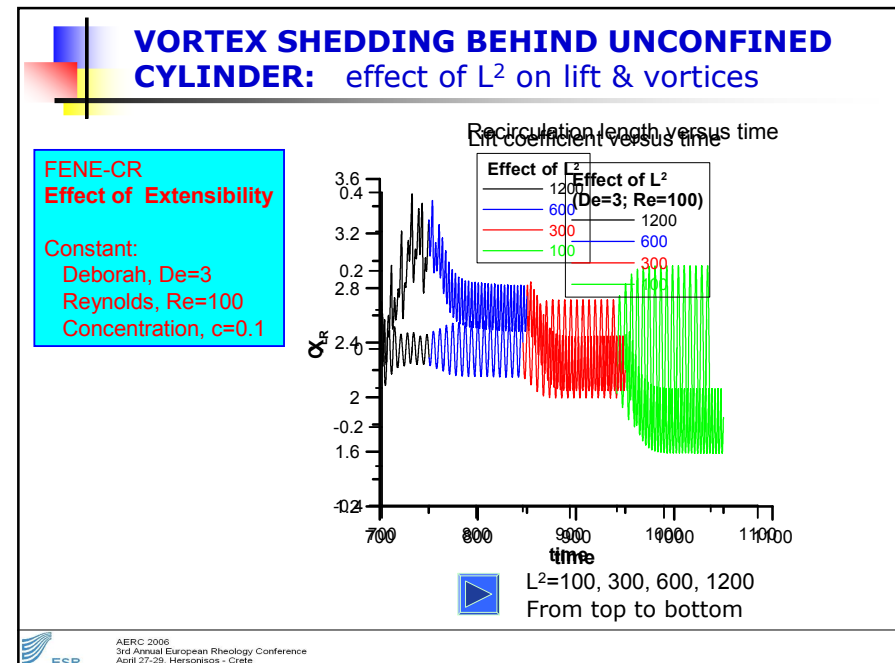
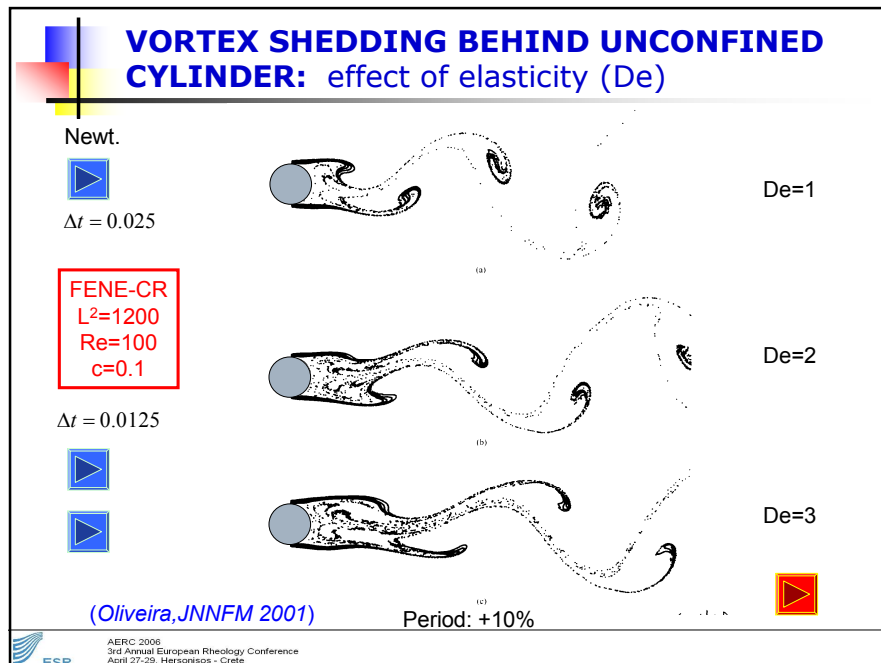


Decay of Strouhal number

Data: Usui et al 1980.

$$St = f \cdot d / U$$

De up, freq. down



VORTEX SHEDDING BEHIND UNCONFINED CYLINDER:

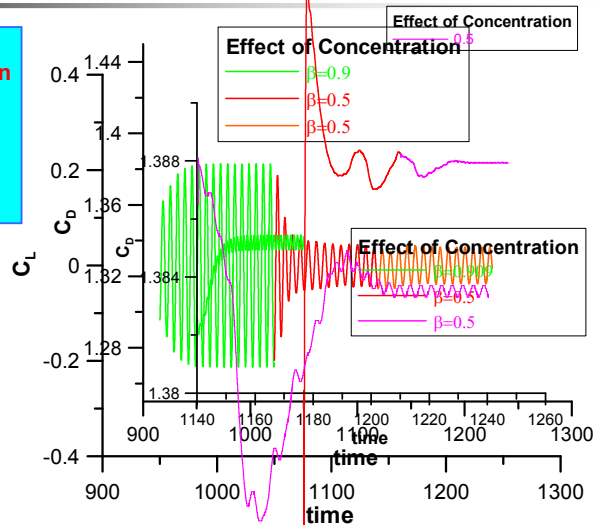
FENE-CR Effect of Concentration

Constant:
Deborah, $De=3$
Reynolds, $Re=100$
Extensibility, $L^2=100$

Diffracted light when
where concentration
rises from 0 to 1 to
 $c=1$

$$\beta = \frac{\eta_s}{\eta_0} = \frac{1}{1+c}$$

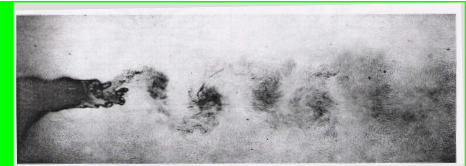
$c=0.1; 1.0; 2.3$
 $c = \frac{\eta_p}{\eta_s} = \frac{1}{\beta} - 1$
 $c=4.0$



VORTEX SHEDDING BEHIND UNCONFINED CYLINDER:

Visualisations in Boger and Walters (by Strauss & Dohmann, TSF 1991):
 $Re=5000$

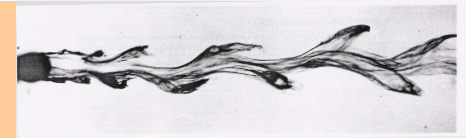
Water:



50 ppm PAC

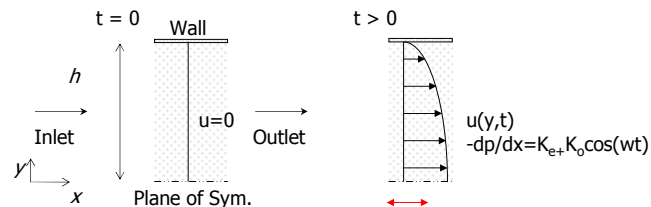


1000 ppm
surfactant



UCM FLOW IN CHANNEL DUE TO PULSATING PRESSURE-GRADIENT (1):

- “Simple” 1D unsteady flow
- Analytical solution available
- Very convenient to assess and develop methods
- Relevance to blood flow
- More complicate than “startup flow”



UCM FLOW IN CHANNEL DUE TO PULSATING PRESSURE-GRADIENT (2):

Elasticity number, $E=1$
Womersley number, $\alpha=4.9$

$$\alpha = h \sqrt{\frac{\omega}{\eta_o / \rho}}$$









Comparison with theory

Simulation results for a period

Line: simulations

Symbols: theory

FUTURE WORK:

-  Near-wall treatment 
-  Devise a “log-law” for stresses
-  Evaluation of gradients
-  Use Log(stress) (Fattal & Kupferman, 2005)
-  3D Calculations at high De 
-  Implement compressibility (unsteady and moving mesh) (sugges. Mackley) (Webster et al)

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